

Progress Report on VICI

Kuniyasu Saitoh

A. Description of my project

Jamming, granular flow, instability of granular flow and dynamic heterogeneities in driven granular systems are my project.

B. Recent talk

(Please see the attached PDF file)

C. Recent papers

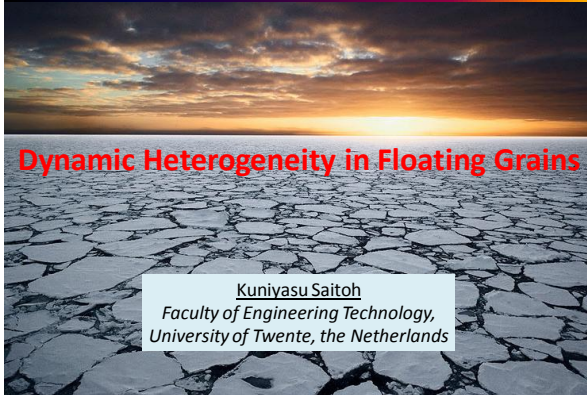
Kuniyasu Saitoh and Hisao Hayakawa, “**Quantitative test of the time dependent Ginzburg-Landau equation for sheared granular flow in two dimension**”, Phys. Fluids (2012), in press.

We have derived the TDGL equation from the hydrodynamic equations of granular gases and adopted to the shear flow. We compare the numerical solutions of the TDGL equation and the hydrodynamic fields, where both results quantitatively same.

Kuniyasu Saitoh, Vanessa Magnanimo, and Stefan Luding, “**Slow dynamics near jamming**”, [AIP Conf. Proc. 1501 \(2012\) 1038](#), *28th International Symposium on Rarefied Gas Dynamics*.

We have studied the jamming transition and critical scaling, especially, the dependence on the size ratio.

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Dynamic Heterogeneity in Floating Grains

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Introduction – Dynamic heterogeneity

Supercooled Liquids

Control parameter

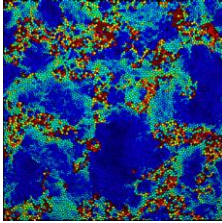
- ❖ *Temperature*

Time scale

- ❖ Self-intermediate scattering function
- ❖ Structural relaxation (rearrangement)
- ❖ Divergence near the glass transition

Length scale

- ❖ Microstructure of glass is almost same with that of liquid (The *static* correlation function does not work)
- ❖ The 4-point *dynamic* correlation function
- ❖ Heterogeneous structure of displacements
- ❖ Divergence near the glass transition temperature



Garrahan (2011) PNAS

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Introduction – Dynamic heterogeneity

Dynamic heterogeneity is universal in disordered systems!

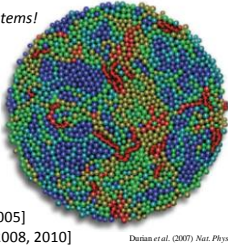
Driven Granular Systems

Control parameters

- ❖ Driving force
- ❖ Density

Examples

- ❖ Air-fluidized grains [Keys, et.al. 2007]
- ❖ Sheared granular systems [Dauchot, et.al. 2005]
- ❖ Horizontally vibrated grains [Lechenault, F. 2008, 2010]



Duran et al. (2007) Nat. Phys.

Our Aim

Understanding of dynamic heterogeneity in more complex systems.

Floating grains
Driven by standing wave, cohesive forces, flow with convection

Experiment

Setup

far field light source

high speed camera

grains

the wavelength = R

air

water

shaker

30 mm

40 mm

Grains
Polystyrene beads distributed on water (mean diameter $\sigma=0.62\text{mm}$, polydisperse)

Driving force
Standing wave generated by a shaker

Control parameter
Area fraction of grains ϕ

Interaction between grains
Capillary force \rightarrow "cohesive"

Data
Trajectories projected on 2-dimensions

Units
time = sec. length = σ

Coarse Graining Method

Trajectories of grains

$$\vec{x}_i(t) = \vec{r}_i(t) + \int_0^t \vec{u}(\vec{x}_i(s), s) ds \quad (i = 1, \dots, N)$$
 Fluctuation Transport by convection

Coarse Graining (CG)

Velocity field

$$\vec{u}(\vec{x}, t) \equiv \frac{\sum \vec{v}_i(t) \varphi_d(\vec{x} - \vec{x}_i(t))}{\sum \varphi_d(\vec{x} - \vec{x}_i(t))}$$

CG function

$$\varphi_d(\vec{x}) = e^{-\frac{(|\vec{x}|/d)^2}{2}} \text{ or } \begin{cases} 1 & (|\vec{x}| \leq d) \\ 0 & (|\vec{x}| > d) \end{cases}$$

Diffusion

Mean Square Displacement

$\vec{x}_i(t)$: Convection is dominant & almost ballistic.
 $\vec{r}_i(t)$: Crossover from sub-diffusion ($C_2 \tau^\beta$) to diffusion ($C_1 \tau^\alpha$).

Crossover time τ_c defined as

$$C_1 \tau_c^\alpha = C_2 \tau_c^\beta \equiv \Delta_{\text{plateau}}$$

diverges near the estimated jamming point $\phi_c = 0.82$ as

$$\tau_c \sim \Delta \phi^{-3.9} \quad (\Delta \phi \equiv \phi_c - \phi)$$

Order Parameter

Displacement
 $r_i(t, \tau) \equiv |\vec{X}_i(t + \tau) - \vec{X}_i(t)|$ where $\vec{X}_i(t) = \vec{x}_i(t)$ or $\vec{r}_i(t)$

Order Parameter

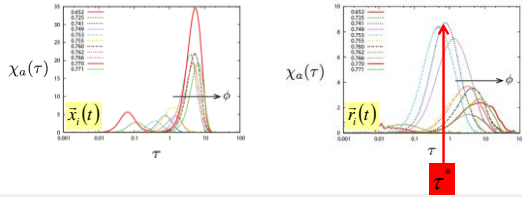
Overlap function
 $w_a(r) = e^{-\frac{r^2}{a^2}}$ or $\begin{cases} 1 & (r \leq a) \\ 0 & (r > a) \end{cases}$

Mobility
 $q_a(t, \tau) \equiv \frac{1}{N} \sum_{i=1}^N w_a(r_i(t, \tau))$

Order parameter
 $Q_a(\tau) = \langle q_a(t, \tau) \rangle_t$

Dynamic susceptibility

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Dynamic Susceptibility

$$\chi_a(\tau) \equiv N \left\langle \left(q_a(t, \tau)^2 \right)_i - \left(q_a(t, \tau) \right)_i^2 \right\rangle$$

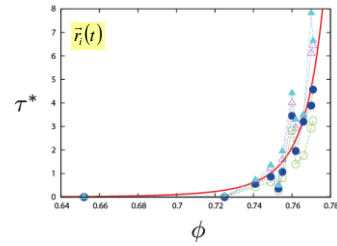
is defined as a *deviation* of mobility & has a *single peak* at τ^* .

$\bar{x}_i(t) : \tau^*$ is not monotonous.

$\bar{r}_i(t) : \tau^*$ monotonously depends on ϕ .

Time Scale

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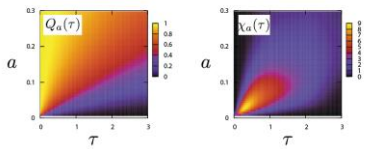
Time Scale

Characteristic time τ^* given by $\bar{r}_i(t)$ diverges near the jamming point as

$$\tau^* \sim \Delta\phi^{-4.2}$$

Width of Overlap Function

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Criteria for width "a"

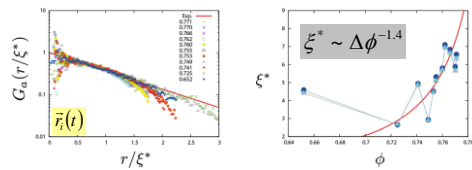
- $a > \sqrt{\Delta_{plateau}}$
- The maximum amplitude of $\chi_a(\tau^*)$.

CG func.	Overlap func.	a (width of OF)
None	Step	σ
	Gauss	σ
Step	Step	0.046σ
($d = 1.6\sigma$)	Gauss	0.042σ
Gauss	Step	0.042σ
($d = \sigma$)	Gauss	0.038σ

Cf.) Horizontally vibrated grains (Lechenault & Dauchot et al.) $a \sim 10^{-2}\sigma$

4-point correlation function

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Definition

$$\chi_a(\tau^*) \equiv 2\pi \int_0^\infty r g_a(r, \tau^*) dr \quad \text{Number density}$$

$$g_a(r, \tau) = \frac{1}{2\pi N} \left\langle \sum_{i,j} \delta(r - |\bar{r}_i(t) - \bar{r}_j(t)|) w_a(r_i(t, \tau)) w_a(r_j(t, \tau)) \right\rangle - \rho \langle q_a(t, \tau) \rangle_i^2$$

Scaling

$$g_a(r, \tau^*) = cr^{-\beta} e^{-r/\xi^*} \Rightarrow G_a(r/\xi^*) = c^{-1} r^\beta g_a(r, \tau^*) \quad (\beta \approx 0.01)$$

Cf.) Horizontally vibrated grains (Lechenault & Dauchot et al.) $\beta \ll 0.15$

Dynamic criticality

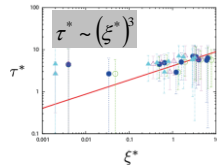
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Dynamic Criticality of Floating Grains

$$\tau^* \sim \Delta\phi^{-\lambda}$$

$$\xi^* \sim \Delta\phi^{-\mu}$$

$$\tau^* \sim (\xi^*)^{\lambda/\mu}$$

Exponents for Driven Granular Systems

	λ	μ	λ/μ	Dimension
Horizontally Vibrated Grains	0.33	0.50	0.65	Quasi-2D
Air-fluidized Grains	1.0	1.7	0.6	Quasi-2D
Sheared Colloids	4	4/3	3	Quasi-2D
Floating Grains	4.2	1.4	3	Quasi-2D

Summary

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Dynamic Heterogeneity in Floating Grains

- ❖ *Floating grains* are driven by *standing wave* & flow with *convection*
- ❖ *CG method* successfully subtracts the additional displacements by *convection*
- ❖ *Crossover time* diverges near *the jamming point*
- ❖ *Time scale* given by susceptibility also diverges near *the jamming point*
- ❖ *Dynamic correlation length* diverges near *the jamming point*
- ❖ *Exponents* obtained from our analysis are consistent with previous works
- ❖ *Our results do not* depend on the choose of CG & overlap functions

Outlook

- ❖ The maximum susceptibility $\chi_a(\tau^*)$
- ❖ Self-intermediate scattering function
- ❖ Self-van Hove function
- ❖ Effect of convection

