Initial Structure and Growth Dynamics of YBa$_2$Cu$_3$O$_{7-\delta}$ during Pulsed Laser Deposition

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The initial heteroepitaxial growth of YBa$_2$Cu$_3$O$_{7-\delta}$ films on SrTiO$_3$(001) substrates during pulsed laser deposition shows a growth-mode transition and a change of growth unit. The growth starts with two blocks, each two-thirds the size of the complete unit cell. The first of these blocks grows in a step-flow fashion, whereas the second grows in the layer-by-layer mode. Subsequent deposition occurs layer-by-layer of complete unit cells. These results suggest that the surface diffusion in the heteroepitaxial case is strongly influenced by the competition with formation energies, which is important for the fabrication of heteroepitaxial devices on the unit cell scale.

Nanoscale structures are undoubtedly a major aspect of contemporary materials science. By attempting to control the shape, crystallinity, and defects of structures at the atomic level, the limits of performance, characterization, and theoretical knowledge are challenged. The high-$T_c$ superconductor YBa$_2$Cu$_3$O$_7$ (YBCO) is a good example of the aforementioned issues. In thin film form, it can be integrated in devices used in electronics, biophysics, and magnetic and x-ray sensors [1]. By nanostructuring of thin films of YBCO it is possible to study the symmetry of the superconducting order parameter, thereby giving fundamental insights into high-$T_c$ superconductivity [2].

Pulsed laser deposition (PLD) is a technique well suited for the fabrication of thin perovskitelike films. The dynamics of the growth during PLD have been extensively studied using reflection high energy electron diffraction (RHEED) [3,4]. However, with RHEED it is not possible to obtain detailed information about the atomic structure during growth, which motivated the use of x rays [5–11]. Here we report on x-ray diffraction measurements during the PLD growth of YBCO on SrTiO$_3$(001) (STO).

The experiments were carried out using a specially designed PLD chamber [12] at the European Synchrotron Radiation Facility on BM26 (DUBBLE) [13], using an x-ray wavelength of 0.775(1) Å. A KrF excimer laser beam (wavelength $\lambda = 248$ nm) was focused on the sintered YBCO target, resulting in a laser fluence of approximately 1.5 J cm$^{-2}$. During deposition the TiO$_2$-terminated substrate temperature was kept at 1053 K, while the oxygen pressure was 0.1 mbar. Five TiO$_2$-terminated substrates, labeled I through V hereafter and differing only in their miscut angles, were used in the measurements.

In order to extract the structure evolution of the film during deposition, a distributed growth model [15] is used. The coverage $\theta$ of the $j$th layer is evaluated as a function of time by solving the following set of coupled differential equations:

$$\frac{d\theta}{dt} = (1 - \lambda_{j-1}) \frac{1}{\tau} (\theta_{j-1} - \theta_j) + \lambda_j \frac{1}{\tau} (\theta_j - \theta_{j+1})$$

with $\tau$ the time needed to deposit a complete monolayer. The interlayer mass transport, $\lambda$, is made coverage-dependent by using the following relation [16]:

$$\lambda_j = 1 - \exp\left(-\alpha \frac{\sqrt{\frac{1}{2} - |\theta_{j-1} - \frac{1}{2}|}}{\theta_{j-1}}\right)$$

with $\alpha$ a dimensionless parameter. When $\alpha \to \infty$, there is perfect layer-by-layer or step-flow growth [17]; for each layer $\lambda_j = 1$, which means that all the material that arrives at a certain level will diffuse to the level below. The other extreme case, $\alpha = 0$, results in Poisson growth. The difference with the model of Ref. [16] is that here the term accounting for a gradual decrease of interlayer mass transport is omitted.

The results of monitoring the $(0,0,0.175)$ reciprocal space point in STO units [18] during deposition of sample I are shown in Fig. 1(a). The oscillating intensity indicates that the surface cycles between smooth and rough morphologies. However, the time that elapses between the consecutive intensity maxima is not constant. The two very first smooth layers (B and C) are each deposited in 2 consecutive intensity maxima. However, the time that elapses between the two very first smooth layers (B and C) is each deposited in 2 consecutive intensity maxima. Subsequent deposition occurs layer-by-layer of complete unit cells. These results suggest that the surface diffusion in the heteroepitaxial case is strongly influenced by the competition with formation energies, which is important for the fabrication of heteroepitaxial devices on the unit cell scale.
where $s$ is a scale factor, $F_{\text{CTR}}$ is the crystal truncation rod (CTR) scattering of the substrate [19], $F_0$ describes the scattering from each of the two first layers, $F_{\text{ybco}}$ is the structure factor of a complete YBCO unit cell taken from the bulk structure [20], and $\Delta z$ is the displacement of the complete unit cell layers YBCO with respect to the layers underneath. For the complete film, a $c$ axis of 11.91(5) Å has been measured, corresponding to the YBCO bulk value at 1053 K. Furthermore, the very first layer is assumed to grow in a SF fashion, which requires one to consider the sum of the resulting intensities of the covered and uncovered substrates as a function of time. From the second layer onward, growth proceeds in the LBL mode. This finding will be discussed in more detail later on. The results of fitting Eq. (3) to the data are listed in Table I and shown in Fig. 1(a).

In order to understand the structure, and thereby $F_0$, of an ultrathin YBCO film, the specular CTR of sample III was measured after deposition of half a unit cell layer. The data [see Fig. 1(b)] were integrated, corrected, and used in a refinement procedure using the ANA-ROD package [21]. The structural model that fits the data best emerges from the following assumptions and restraints, which were derived from previously published results on the growth and (surface) structure of YBCO and related materials. The
model consists of perovskitelike blocks \( ABX_3 \), with \( A = Y_{1/3}Ba_{2/3}, B = Cu, \) and \( X = O_{2/3} \). These represent the average structure in the case that the \( Y \) and \( Ba \) atoms would not be ordered in the bulk YBCO structure. The stoichiometry is fixed in the refinement procedure, where the best fit is obtained when half of the oxygen atoms are missing in the surface copper-oxide planes [Cu2-O6 and half of the Cu1-O4 plane] [22]. The atomic stacking across the interface is assumed to continue the network of oxygen octahedra, with Ti at their centers in the substrate and Cu in the film. Starting values for the atomic positions in the refinement were derived from the bulk structures \( La_{2}Y_{2}Cu_{2}Ti_{2}O_{11} \) [23] and YBCO [20]. The results are listed in Table II, and a schematic picture of the structure is shown in Fig. 1(c).

The best fit model consists of two layers of the \( ABX_3 \) blocks, denoted 1 and 2 in Fig. 1(c). The layer at the interface (1) is fully covering the substrate and the one on top (2) is covering half of the surface. Such a morphology would give rise to step-heights on the surface of about \( 4 \) and possibly \( 8 \) Å, which is in good agreement with previous studies on ultrathin RBCO films, with \( R = Y \) [24], Sm [25]. Test models having different morphologies are clearly distinguished as shown in Fig. 1(b). This result suggests that in the case of sample III the nucleation proceeds differently compared to sample I, which is, as discussed later on, due to the different miscut angles of the substrates. However, using the structural model to calculate the structure factor in the case of fully occupied layers 1 and 2 at point \( l = 0.175 \) results in \( |F_0| = 159 \) and \( \phi_0 = 1.64 \). This agrees rather well with the results obtained for \( F_0 \) shown in Table I and suggests strongly that the growth unit during initial deposition of sample I resembles the presented structural model very much.

It was not possible to obtain a more extensive data set, since about 1 h after deposition the ultrathin film was found to be unstable. A possible mechanism for the deterioration of the structure is that the copper-oxide planes at the surface are slowly reduced. The oxygen leaves the surface at defects that are reached by in-plane diffusion. When the total oxygen content drops below a critical value, the structure becomes unstable and a, so far unidentified, new phase is formed. It must be emphasized that the current method and data do not allow one to determine and refine all the details of the oxygen content, of which the uncertainty is estimated to be somewhere between 0.5 and 1 oxygen atom. Nevertheless, the perovskitelike model structure provides a framework of heavier elements, which allows it to be refined and which consistently explains all the observations.

Unlike room temperature experimental results for thicker films [26], our best fit model shows that the \( A \) site is randomly occupied by \( Ba \) and \( Y \). This finding is further supported by all the tried test models (including ordering of \( Y \) and \( Ba \) over the \( A \) sites) resulting in approximately equal Ti-Cu1 and Cu1-Cu2 distances. In case of \( A \) site ordering this would not be expected, since an \( YCuO_3 \) block is considerably smaller than a \( BaCuO_3 \) block, 3.42 and 4.18 Å, respectively. The mechanism could be such that with increasing thickness the \( A \) site orders, simply because there is enough material. The ordering then leads to different interfacial stacking sequences substrate-TiO_2-BaCuO_x-BaCuO_y-YCuO_z... and substrate-TiO_2-BaCuO_x-YCuO_y-BaCuO_z... [24,26], thereby forming commonly observed antiphase boundaries [26] and possibly leading to different surface terminations [22].

Figure 1(d) shows different shapes for the first intensity oscillation of samples I and III. For sample I the intensity increases linearly, whereas for sample III it decreases parabolically. These differences indicate SF (sample I) and LBL (sample III) growth. In the former case there is no roughening of the surface and the scattering is given by the sum of intensities of the covered and the uncovered substrates. The latter situation implies nucleation of two-
dimensional islands on the terraces. The resulting scattering arises from a mixed contribution on the scale of the x-ray coherence length [19] of the bare and covered substrates. This leads to the inequality \[ |F_{CTR} + \epsilon F_0|^2 < |F_{CTR}|^2 + (1 - \epsilon) |F_{CTR}|^2 + \epsilon |F_{CTR} + F_0|^2, \] where \( \epsilon \) is an arbitrary number with \( 0 < \epsilon < 1 \), showing that the intensity expected from SF is higher. Since simultaneous nucleation on the terraces and at the step-edges is conceivable, a parameter \( \beta \) is defined determining the fraction of material nucleating by step-flow. The scattered intensity is then given by \( I_{dSF} = \beta I_{dF} (1 - \beta) I_{dLBL} \), with \( I_{SF} \) and \( I_{LBL} \) as in Eq. (4). Figure 1(d) shows simulations for the very first intensity oscillation of sample III for \( \beta = 0.65 \) and 0.00.

The results for sample I lead to the conclusion that the diffusion length changes during growth. In order to estimate the initial diffusion length, \( l_{dinit} \), experiments have been carried out on five TiO\(_2\)-terminated substrates, having different miscut angles with respect to the [001] direction. The results, listed in Table III, show a clear correlation between the miscut angle and the scattered intensity increasing or decreasing upon starting the deposition.

The present results indicate that the diffusion length of the very first layer is around 150 nm, much larger than reported values for YBCO growth ranging from 7 to 22 nm [27,28]. A rough estimate of the difference in the activation energy can be obtained from the diffusion relation by \( \Delta E = \ln(I_d/I_{dinit}) \times 2k_BT = -0.4(1) \text{ eV} \) [29]. This would represent a substantial part of the activation barrier for diffusion being around 1 eV for complex oxide growth at about 1000 K [30]. An explanation might be that when the terminating atomic layer changes from STO to YBCO, the energy landscape for diffusion changes. Furthermore, with initial nucleation taking place only at the substrate steps (SF), the additional energy barrier for diffusion over the step, the so-called Ehrlich-Schwöbel barrier can change. This may result in it being easier for adatoms to ascend a step instead of descend [31]. Indeed, a scanning tunneling microscopy study of ultrathin YBCO films has revealed that nucleation occurs preferentially at the upper step [24], quite different from homoepitaxy.

In the case of heteroepitaxial deposition of SrRuO\(_3\) (SRO) on STO(001), a change in growth mode during initial nucleation has been observed as well [32]. However, in that case growth started as LBL and changed to SF. Interestingly, the lattice mismatch between SRO and STO leads to compressive strain in the film, opposite to the present system. Compressive strain has been found to lower the activation barrier for surface diffusion [33], opposite from tensile strain, and is therefore another important parameter in heteroepitaxial growth.

In conclusion, the change from one atomic structure to another involves dramatic changes in the surface free energies, thereby altering the growth completely. In the case of YBCO growth on STO(001) both the growth unit and the growth mode change during the course of the nucleation of the two first layers. These results are very important in view of the fabrication of real heteroepitaxial devices on the unit cell scale.

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[17] Within this model, it does not contain any information about the island-size distribution. Merely the interlayer mass transport is described. Therefore within this model step-flow and perfect layer-by-layer are indistinguishable.

[18] This is very close to the “anti-Bragg” position of bulk YBCO.


[29] The error represents the scatter of the different literature values.


